# FUNDAMENTAL INEQUALITIES FOR THE BOUNDS ON THE EFFECTIVE TRANSPORT COEFFICIENTS OF TWO-PHASE MEDIA

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Summary By starting from the truncated power expansions given at a number of discrete real points we establish the general inequalities for the effective transport coefficients Q of macroscopically isotropic two phase media. The inequalities obtained provide the new upper and lower bounds on Q the best with respect to the available power series coefficients. For specific cases these new bounds reduce to the classical ones due to Wienner and Hashin-Shtrikman. They are also coincide with the estimations of Bergman [1] and Milton [3] derived for the fitting problem. Many illustrative examples show the usefullness of the results obtained..

#### FORMULATION OF THE PROBLEM

It is well-known that the effective transport coefficient of a two-phase medium with macroscopic isotropic symmetry has a Stieltjes integral representation:

$$f_1(x) = \frac{Q(x) - 1}{x} = \int_0^1 \frac{d\gamma_1(u)}{1 + xu}, \ d\gamma_1(u) > 0, \ x = \frac{q_2}{q_1} - 1, \ Q(x) = \frac{q_{ef}(x)}{q_1}; \ f_1(-1) \le 1.$$
 (1)

Here  $q_{ef}$  denotes the effective modulus of the composite material, while  $q_1, q_2$  describe the physical properties of its constituents. Assume that the truncated power expansions of  $f_1(x)$  at  $x_1, x_2, ..., x_N, -1$  be given, where  $-1 < x_j, j = 1, 2, ..., N$ ,

$$f_1(x) = f_1(-1) + O(x+1); \ f_1(-1) \le 1; \ f_1(x) = \sum_{i=0}^{p_j-1} c_{ij}(x-x_j)^i + O((x-x_j)^{p_j}), \ j = 1, 2, ..., N$$
 (2)

The following abbreviated notation for the truncated power series (2) will also be used:

$$f_1(x) = f_1|_{\boldsymbol{x}}^{\boldsymbol{p}}(x-\boldsymbol{x}) + O(x-\boldsymbol{x})^{\boldsymbol{p}}, \ \boldsymbol{x} = [x_1, x_2, ..., x_N, -1], \ \boldsymbol{p} = [p_1, p_2, ..., p_N, 1], \ f_1(-1) \le 1$$
 (3)

The aim of this contribution is to establish, starting from the power expansions (3), the general inequalities for the Stieltjes function  $f_1(x)$ , and next via (1) the best upper and lower bounds on the effective coefficient Q(x).

#### THE METHOD OF SOLUTION

In the first step the inequality  $f_1(-1) \le 1$  is replaced formally by the equality  $f_1(-1) = 1$ , cf. (3). The input data (3) transform to

$$f_1(x) = f_1|_x^p(x-x) + O(z-x)^p, \ x = [x_1, x_2, ..., x_N, -1], \ p = [p_1, p_2, ..., p_N, 1], \ f_1(-1) = 1.$$
 (4)

Next by introducing  $P_0=0$ ,  $P_j=\sum_{i=1}^j p_i$ , j=1,2,...,N;  $P=P_N+1$  we expand the original Stieltjes function  $f_1(x)$  given by (1) to the new Stieltjes one  $f_P(x)$ 

$$f_{P_0+1}(x) = \frac{f_{P_0+1}(x_1)}{1+(x-x_1)f_{P_0+2}(x)}, \ f_{P_0+2}(x) = \frac{f_{P_0+2}(x_1)}{1+(x-x_1)f_{P_0+3}(x)}, \ \dots, \ f_{P_1}(x) = \frac{f_{P_1}(x_1)}{1+(x-x_1)f_{P_1+1}(x)},$$

$$f_{P_1+1}(x) = \frac{f_{P_1+1}(x_2)}{1+(x-x_2)f_{P_1+2}(x)}, \ f_{P_1+2}(x) = \frac{f_{P_1+2}(x_2)}{1+(x-x_2)f_{P_1+3}(x)}, \dots, \ f_{P_2}(x) = \frac{f_{P_2}(x_2)}{1+(x-x_2)f_{P_2+1}(x)},$$

$$(5)$$

$$f_{P_{N-1}+1}(x) = \frac{f_{P_{N-1}+1,\,}(x_N)}{1+(x-x_{N-1})f_{P_{N-1}+2,\,}(x)},..,\ \ f_{P_N}(x) = \frac{f_{P_N}(x_N)}{1+(x-x_N)f_P(x)},\ \ f_P(x) = f_P(-1) + O(x+1).$$

It is convenient to rewrite (5) as follows

$$f_1(x) = f_1_x^p(x, f_P(x)) = f_1_{x_1, x_2, \dots, x_N, 1}^{p_1, p_2, \dots, p_N, 1}(x, f_P(x)), \text{ where } f_P(x) = \int_0^1 \frac{d\gamma_P(u)}{1+xu} \text{ and } f_P(-1) = w_p.$$
 (6)

For any values of  $f_P(x) \in [0, w_p]$ ,  $x \in [-1, \infty]$ , the continued fraction expansion  $f_{1x}(x, f_P(x))$  satisfies the input data (4). Let us introduce now the function  $F_P(x, u)$  estimating  $f_1(x)$  in the following way

$$f_1(x) \in \{F_P(x, u); -1 \le u \le 0\}, \text{ where } F_P(x, u) = f_1^{p_1, p_2, \dots, p_N, 1}_{x_1, x_2, \dots, x_N, 1}(x, w_P(1+u)),$$
 (7)

The sign of the first derivative of  $F_P(x, u)$  with respect to u, obtained after lengthy calculations is equal to:

$$\operatorname{sign}\left(\frac{\partial F_P(x,u)}{\partial u}\right) = (-1)^0, = (-1)^{P_1}, \dots, = (-1)^{P_N} \text{ if } -1 < x < x_1, \ x_1 < x < x_2, \dots, \ x_N < x < \infty,$$
 (8)

where  $P_j = \sum_{i=1}^j p_i$ , j = 1, 2, ..., N;  $P = P_N + 1$ . From (8) the inequalities follow:

$$(-1)^{0} F_{P}(x,-1) \leq (-1)^{1} f_{1}(x) \leq (-1)^{1} F_{P}(x,0) \quad \text{if} \quad -1 < x \leq x_{1},$$

$$(-1)^{P_{1}} F_{P}(x,-1) \leq (-1)^{P_{1}+1} f_{1}(x) \leq (-1)^{P_{1}+1} F_{P}(x,0) \quad \text{if} \quad x_{1} \leq x \leq x_{2},$$

$$(-1)^{P_{N}} F_{P}(x,-1) \leq (-1)^{P_{N}+1} f_{1}(x) \leq (-1)^{P_{N}+1} F_{P}(x,0) \quad \text{if} \quad x_{N} \leq x < \infty.$$

$$(9)$$

The functions  $F_P(x,0)$  and  $F_P(x,-1)$  are the multipoint Padé approximant  $[m_P/n_P](x)$  and  $[m_{P-1}/n_{P-1}](x)$  to  $f_1(x)$ , where  $[m_P/n_P] = \frac{a_0 + a_1 x^1 + \dots + a_{m_P} x^{m_P}}{1 + b_1 x^1 + \dots + b_{n_P} x^{n_P}}$ ,  $m_P = P - 1 - n_P$ ,  $n_P = E(P/2)$ ,  $P = \sum_{j=1}^N p_j + 1$ . Relations (9) provide the fundamental inequalities for the multipoint Padé approximants to the Stieltjes function  $f_1(x)$  representing via (1) the effective transport coefficient Q(x).

# FUNDAMENTAL INEQUALITIES

Let  $L_R(x) = \sum_{j=1}^N p_j H(x-x_j)$  determining the total number  $p_1 + p_2 + ...p_s$  of the coefficients of the power expansions of  $f_1(x)$  available at points  $x_1, x_2, ..., x_s \leq x$  be given, see (4). By introducing the piecewise continuous function  $M_R(x) = (L_R(x) \text{ if } -\infty < x < 0 \text{ or } L_R(x) + 1 \text{ if } -\infty < x < 0)$ , the relations (9) can be easily transformed to the following inequalities for the multipoint Padé approximants  $x[m_R/n_R]$  and  $x[m_{R+1}/n_{R+1}]$  to the effective transport coefficient Q(x) - 1, cf. (1),

$$(-1)^{M_R(x)} x [m_{R+1}/n_{R+1}](x) \le (-1)^{M_R(x)} (Q(x) - 1) \le (-1)^{M_R(x)} x [m_R/n_R](x), \ x \in [-1, \infty).$$

$$= \frac{20}{15} \begin{bmatrix} & & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & \\ & & & & \\ & & &$$

Fig. 1: Three- and four-point Padé approximants to the function  $Q(x) - 1 = xf_1(x)$ ,  $f_1(x) = \frac{\ln(0.5(2+x))}{x}$ . representing the bounds on Q(x) - 1 predicted by the fundamental inequality (10). Here  $x_1 = -1$ ,  $x_2 = 0$ ,  $x_3 = 999$ ,  $x_4 = 999999$ .

Inequalities (10) provides the best upper and lower bounds on Q(x) - 1 as a function of  $L_R(x)$  depending on the given numbers of coefficients of power series (4). The multipoint Padé approximant bounds given by (10) generalize all previous bounds reported in the literature, cf. [1,2,3,4]. From that point of view the general bounds (10) are new.

## PARTICULAR CASES

From Q(x) = 1 + O(x), Q(x) = Q(-1) + O(x+1),  $Q(-1) \le 1$  the elementary bounds follow.

$$(-1)^{H(x)}(1+x) \le (-1)^{H(x)}Q(x) \le (-1)^{H(x)}, \ x \in [-1, \infty), \tag{11}$$

For  $Q(x) = 1 + \varphi_2 x + O(x^2)$ , Q(x) = Q(-1) + O(x+1),  $Q(-1) \le 1$  the Wienner bounds result.

$$(-1)^{2H(x)} \left( 1 + \frac{\varphi_2 x}{1 + \varphi_1 x} \right) \le (-1)^{2H(x)} Q(x) \le (-1)^{2H(x)} \left( 1 + \varphi_2 x \right), \ x \in [-1, \infty). \tag{12}$$

For  $Q(x) = 1 + \varphi_2 x + 0.5 \varphi_2 \varphi_1 x^2 + O(x^3)$ , Q(x) = Q(-1) + O(x+1),  $Q(-1) \le 1$  the H–S bounds are obtained as  $Q(x) = 1 + \varphi_2 x + 0.5 \varphi_2 \varphi_1 x^2 + O(x^3)$ .

$$(-1)^{3H(x)} \left( 1 + \frac{\varphi_2 x + 0.5\varphi_2 x^2}{1 + 0.5(1 + \varphi_1)x} \right) \le (-1)^{3H(x)} Q(x) \le (-1)^{3H(x)} \left( 1 + \frac{\varphi_2 x}{1 + 0.5\varphi_1 x} \right), \ x \in [-1, \infty).$$
 (13)

Here  $\varphi_1$  and  $\varphi_2$  denote the volume fractions of the first and second component of the two-phase composite.

## CONCLUSIONS

By using special multipoint continued fraction technique, from several truncated power series we have derived the general inequalities for the multipoint Padé approximants to the effective transport coefficients ( dielectric or diffusion constants, magnetic permeabilities, thermal or electrical conductivities) of two-phase media. Those inequalities are new and provide the best upper and lower bounds on Q(x) over the entire class of rational functions. Note that they are obtained in a unified and coherent form as a function of  $M_R(x)$  depending on given numbers of coefficients of the power expansions of Q(x) only. Moreover, they generalize all previously known relevant bounds, cf. [1,2,3,4].

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