

$$\begin{aligned} (-1)^0 F_P(x, -1) &\leq (-1)^1 f_1(x) \leq (-1)^1 F_P(x, 0) && \text{if } -1 < x \leq x_1, \\ (-1)^{P_1} F_P(x, -1) &\leq (-1)^{P_1+1} f_1(x) \leq (-1)^{P_1+1} F_P(x, 0) && \text{if } x_1 \leq x \leq x_2, \\ \vdots & && \vdots \\ (-1)^{P_N} F_P(x, -1) &\leq (-1)^{P_N+1} f_1(x) \leq (-1)^{P_N+1} F_P(x, 0) && \text{if } x_N \leq x < \infty. \end{aligned} \quad (9)$$

The functions $F_P(x, 0)$ and $F_P(x, -1)$ are the multipoint Padé approximant $[m_P/n_P](x)$ and $[m_{P-1}/n_{P-1}](x)$ to $f_1(x)$, where $[m_P/n_P] = \frac{a_0 + a_1 x^1 + \dots + a_{m_P} x^{m_P}}{1 + b_1 x^1 + \dots + b_{n_P} x^{n_P}}$, $m_P = P - 1 - n_P$, $n_P = E(P/2)$, $P = \sum_{j=1}^N p_j + 1$. Relations (9) provide the fundamental inequalities for the multipoint Padé approximants to the Stieltjes function $f_1(x)$ representing via (1) the effective transport coefficient $Q(x)$.

FUNDAMENTAL INEQUALITIES

Let $L_R(x) = \sum_{j=1}^N p_j H(x - x_j)$ determining the total number $p_1 + p_2 + \dots + p_s$ of the coefficients of the power expansions of $f_1(x)$ available at points $x_1, x_2, \dots, x_s \leq x$ be given, see (4). By introducing the piecewise continuous function $M_R(x) = (L_R(x) \text{ if } -\infty < x < 0 \text{ or } L_R(x) + 1 \text{ if } -\infty < x < 0)$, the relations (9) can be easily transformed to the following inequalities for the multipoint Padé approximants $x[m_R/n_R]$ and $x[m_{R+1}/n_{R+1}]$ to the effective transport coefficient $Q(x) - 1$, cf. (1),

$$(-1)^{M_R(x)} x[m_{R+1}/n_{R+1}](x) \leq (-1)^{M_R(x)} (Q(x) - 1) \leq (-1)^{M_R(x)} x[m_R/n_R](x), \quad x \in [-1, \infty). \quad (10)$$

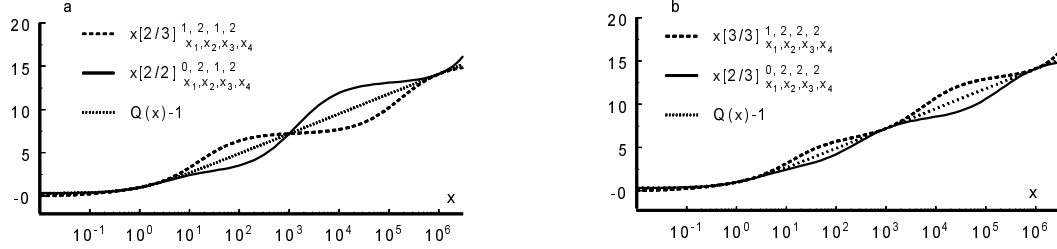


Fig. 1: Three- and four-point Padé approximants to the function $Q(x) - 1 = x f_1(x)$, $f_1(x) = \frac{\ln(0.5(2+x))}{x}$, representing the bounds on $Q(x) - 1$ predicted by the fundamental inequality (10). Here $x_1 = -1$, $x_2 = 0$, $x_3 = 999$, $x_4 = 999999$.

Inequalities (10) provides the best upper and lower bounds on $Q(x) - 1$ as a function of $L_R(x)$ depending on the given numbers of coefficients of power series (4). The multipoint Padé approximant bounds given by (10) generalize all previous bounds reported in the literature, cf. [1,2,3,4]. From that point of view the general bounds (10) are new.

PARTICULAR CASES

From $Q(x) = 1 + O(x)$, $Q(x) = Q(-1) + O(x + 1)$, $Q(-1) \leq 1$ the elementary bounds follow.

$$(-1)^{H(x)} (1 + x) \leq (-1)^{H(x)} Q(x) \leq (-1)^{H(x)}, \quad x \in [-1, \infty), \quad (11)$$

For $Q(x) = 1 + \varphi_2 x + O(x^2)$, $Q(x) = Q(-1) + O(x + 1)$, $Q(-1) \leq 1$ the Wiener bounds result.

$$(-1)^{2H(x)} \left(1 + \frac{\varphi_2 x}{1 + \varphi_1 x}\right) \leq (-1)^{2H(x)} Q(x) \leq (-1)^{2H(x)} (1 + \varphi_2 x), \quad x \in [-1, \infty). \quad (12)$$

For $Q(x) = 1 + \varphi_2 x + 0.5\varphi_2\varphi_1 x^2 + O(x^3)$, $Q(x) = Q(-1) + O(x + 1)$, $Q(-1) \leq 1$ the H-S bounds are obtained

$$(-1)^{3H(x)} \left(1 + \frac{\varphi_2 x + 0.5\varphi_2\varphi_1 x^2}{1 + 0.5(1 + \varphi_1)x}\right) \leq (-1)^{3H(x)} Q(x) \leq (-1)^{3H(x)} \left(1 + \frac{\varphi_2 x}{1 + 0.5\varphi_1 x}\right), \quad x \in [-1, \infty). \quad (13)$$

Here φ_1 and φ_2 denote the volume fractions of the first and second component of the two-phase composite.

CONCLUSIONS

By using special multipoint continued fraction technique, from several truncated power series we have derived the general inequalities for the multipoint Padé approximants to the effective transport coefficients (dielectric or diffusion constants, magnetic permeabilities, thermal or electrical conductivities) of two-phase media. Those inequalities are new and provide the best upper and lower bounds on $Q(x)$ over the entire class of rational functions. Note that they are obtained in a unified and coherent form as a function of $M_R(x)$ depending on given numbers of coefficients of the power expansions of $Q(x)$ only. Moreover, they generalize all previously known relevant bounds, cf. [1,2,3,4].

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